

Spin foam renormalization à la GFT: status and prospects

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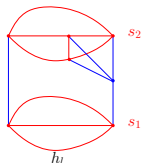
Loops 2017 – Warsaw – July 4th 2017

- 1 Renormalization of spin foams: why and how?
- 2 A non-trivial but tractable GFT theory space
- 3 Perturbative and non-perturbative results
- 4 Conclusion and outlook

- Loop Quantum Gravity proposes **kinematical states** describing (spatial) quantum geometry [Ashtekar, Rovelli, Smolin, Lewandowski... '90s; Dittrich, Geiller, Bahr '15].

Loop Quantum Gravity and Spin Foams

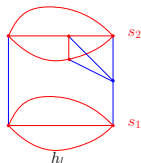
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→ amplitudes $\mathcal{A}_{s,C}$ associated to a 2-complex C with boundary spin-network state s .



→ $\mathcal{A}_{s,C}(h_l)$ with $s = s_1 \cup s_2$
(generalized lattice gauge theory amplitude)

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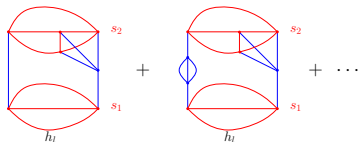
- Two *a priori* distinct strategies to extract \mathcal{A}_s , and therefore **complete the definition of spin foam models**:
 - 1 **Lattice interpretation**: refining and coarse-graining C (and s)
⇒ $\mathcal{A}_s \equiv \lim_{C \rightarrow \infty} \mathcal{A}_{s,C}$ [Dittrich, Bahr, Steinhaus, Martin-Benito, Delcamp,... '10s]
 - 2 **QFT interpretation**: quantum space-time histories, to be summed over
⇒ $\mathcal{A}_s \equiv \sum_{C|\partial C=s} w_C \mathcal{A}_{s,C}$ [De Pietri, Rovelli, Freidel, Oriti... '00s, '10s]

QFT interpretation:

Superposition principle \rightarrow spin foams are viewed as Feynman amplitudes, to be summed over:

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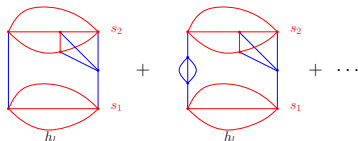


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This may be implemented by a **Group Field Theory**:

[Talk by Oriti]

- a (Euclidean) **quantum field theory** for a (complex) field φ ;
- defined on d -copies of a **local symmetry group** of gravity: $\varphi = \varphi(g_1, \dots, g_d)$;
- with **combinatorially non-local** interactions.

\Rightarrow This embedding of spin foam amplitudes into a GFT uniquely fixes the combinatorial weights w_C .

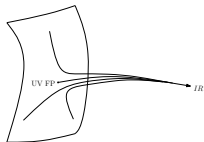
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Physical motivations:

1 **flow** of the quantum dynamics under coarse-graining?



2 **quantization / discretization ambiguities**: what are the **universal features** of the known models?
[EPRL, DL, BO, ...]

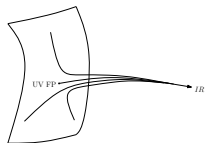
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GFT renormalization programme:

step-by-step generalization of standard renormalization techniques, until we are able to tackle 4d quantum gravity proposals.

[Benedetti, Ben Geloun, Bonzom, SC, Dine Samary, Gurau, Lahoche, Oriti, Rivasseau, Vignes-Tourneret...]

General structure of a GFT and long-term objectives

Typical structure of a GFT: (complex) field $\varphi(g_1, \dots, g_d)$, $g_\ell \in G$, with partition function

$$Z = \int [\mathcal{D}\varphi]_\Lambda \exp \left(-\bar{\varphi} \cdot \mathcal{K} \cdot \varphi + \sum_{\{\mathcal{V}\}} t_{\mathcal{V}} \mathcal{V} \cdot (\bar{\varphi}\varphi)^{p_{\mathcal{V}}} \right) = \sum_{\mathcal{G}} \prod_i (t_{\mathcal{V}_i})^{n_{\mathcal{V}_i}(\mathcal{G})} \mathcal{A}_{\mathcal{G}, \Lambda}$$

Main objectives of the GFT research programme:

- 1 Model building: define the **theory space**.
e.g. spin foam amplitudes $\mathcal{A}_{\mathcal{G}}$ + combinatorial considerations (tensor models) $\rightarrow d, G, \mathcal{K}, \{\mathcal{V}\}$ and $[\mathcal{D}\varphi]_\Lambda$.
- 2 Perturbative definition: prove that the spin foam expansion is **consistent** in some range of Λ .
e.g. perturbative renormalization.
- 3 Systematically explore the theory space: **effective continuum regime reproducing GR** in some limit?
e.g. functional RG, constructive methods, condensate states...

A non-trivial but tractable GFT theory space

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Simple enough but non-trivial set-up: [Ben Geloun, Rivasseau '11; SC, Oriti, Rivasseau '12; ...]

- $G =$ **compact Lie group**: $SU(2)$ for definiteness.
- Standard bosonic kinetic term: $\mathcal{K} = m^2 - \sum_{\ell=1}^d \Delta_\ell$.
- $[\mathcal{D}\varphi]_\Lambda =$ integration over **gauge invariant fields**, with **cut-off Λ on large spins**.
- $\{\mathcal{V}\} =$ **{connected trace invariants}** $\rightarrow \infty$ -dimensional theory space.

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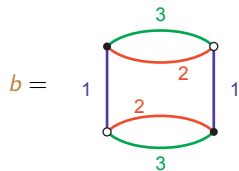
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\Rightarrow sum over $SU(2)$ lattice gauge theory amplitudes.

Trace invariants

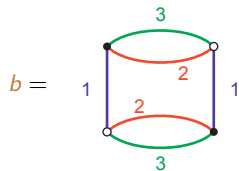
Trace invariants of fields $\varphi(g_1, g_2, \dots, g_d)$ labelled by d -colored graphs / bubbles b :



$$\mathrm{Tr}_b(\varphi, \bar{\varphi}) = \int [dg_i]^6 \bar{\varphi}(g_6, g_2, g_3) \varphi(g_1, g_2, g_3) \\ \varphi(g_6, g_4, g_5) \bar{\varphi}(g_1, g_4, g_5)$$

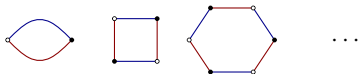
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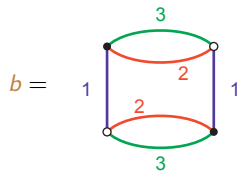
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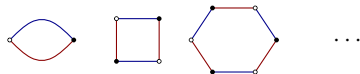
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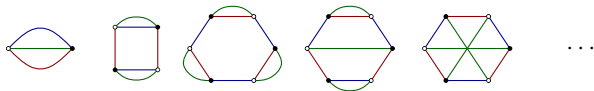


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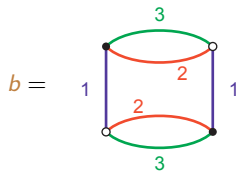


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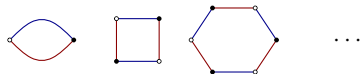
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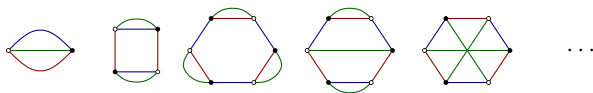


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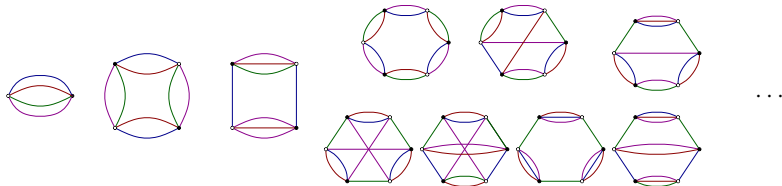
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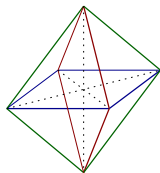
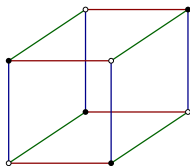
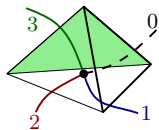
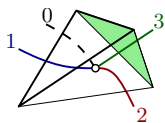
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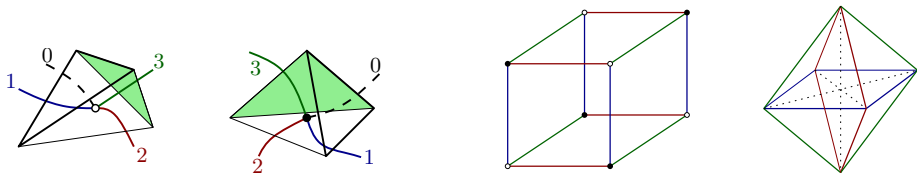
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Colored cell decompositions of pseudo-manifolds

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Theorem: [Pezzana '74] Any PL manifold can be represented by a colored graph. In general, a $(d + 1)$ -colored graph represents a triangulated **pseudo-manifold** of dimension d .

\Rightarrow **Crystallisation theory** [Cagliardi, Ferri et al. '80s]; introduced in GFTs / tensor models in '09 [Gurau '09...].

Nice book and reviews: "Random Tensors", Gurau (2016) and [Gurau, Ryan '11; Ryan '16]

- In order to connect to spin foams, the GFT path integral must be restricted to **gauge invariant fields**:

$$\forall h \in \text{SU}(2), \quad \varphi(g_1 h, g_2 h, \dots, g_d h) = \varphi(g_1, g_2, \dots, g_d)$$

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- Choose the **standard bosonic propagator** $C = \left(m^2 - \sum_{\ell=1}^d \Delta\right)^{-1}$. It is regularized by means of a **cut-off on large spins**:

$$\sum_{\ell=1}^3 j_{\ell}(j_{\ell} + 1) \lesssim \Lambda^2$$

Gauge invariance condition and Feynman amplitudes

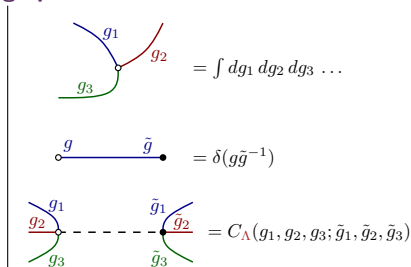
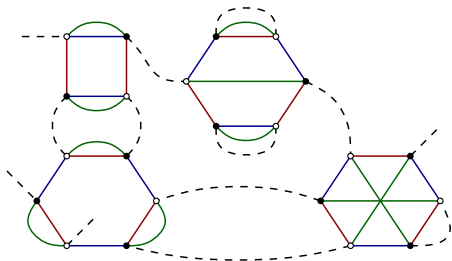
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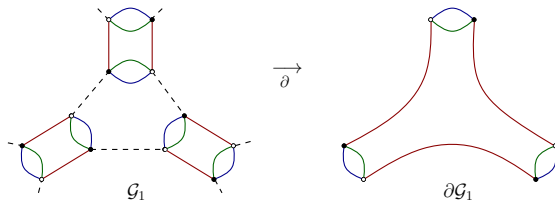
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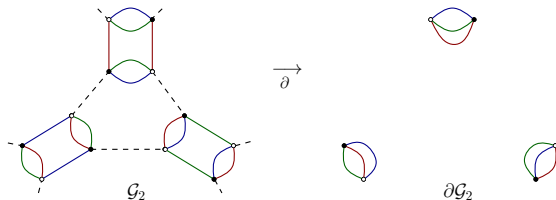
- The Feynman diagrams are **$(d+1)$ -colored graphs**:



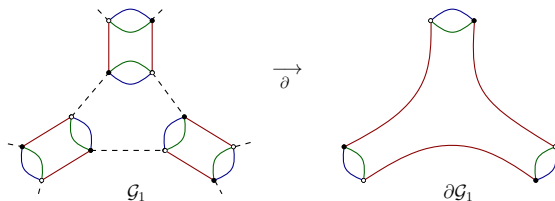
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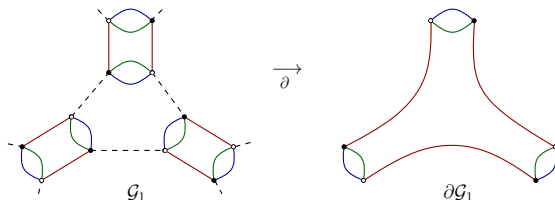
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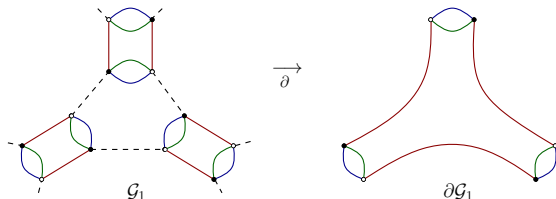


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This is **not automatically possible**, but it has been proven to work in the present theory space. This relies on a **delicate interplay** between the divergent structure of the theory and the topology of the divergent Feynman diagrams:

- the divergences are supported on **flat bulk holonomies**;
- the leading order Feynman diagrams ('melonic graphs') and their boundaries have **trivial topologies**.

[SC, Oriti, Rivasseau '13...]

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- Perturbatively renormalizable actions for $d = 3$ and $d = 4$ $SU(2)$ GFTs:

$$\begin{aligned}
 S_\Lambda(\varphi, \bar{\varphi}) &\stackrel{(d=3)}{=} -\bar{\varphi} \cdot \Delta \cdot \varphi + t_2^\Lambda \text{ (loop) } + t_4^\Lambda \text{ (cylinder) } + t_{6,1}^\Lambda \text{ (hexagon) } + t_{6,2}^\Lambda \text{ (diamond) } \\
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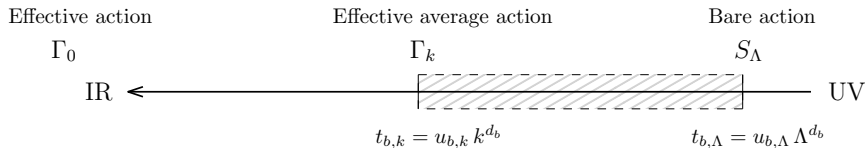
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- Perturbative **beta functions** can be extracted \rightarrow prescription for how to consistently adjust t_b^Λ as one changes Λ . [SC '14...]

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e.g. functional RG, constructive methods, condensate states...
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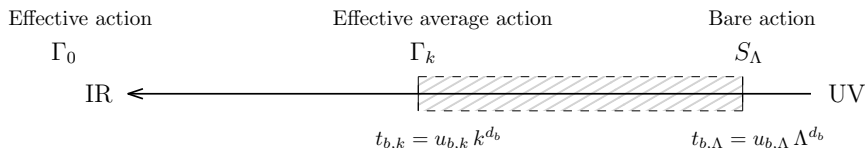
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Effective average action [Wetterich '93, Morris '93...] [Benedetti, Ben Geloun, Oriti '14; Benedetti, Lahoche '15...]



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e.g. functional RG, constructive methods, condensate states...

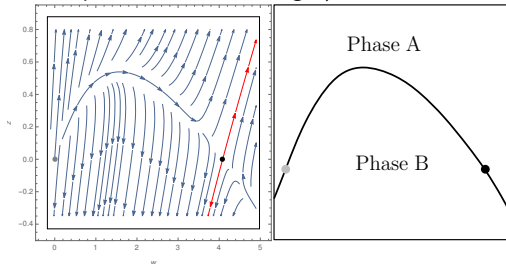
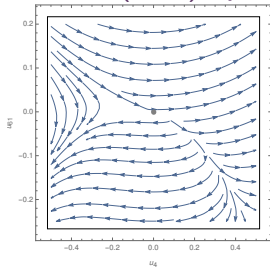
Effective average action [Wetterich '93, Morris '93...] [Benedetti, Ben Geloun, Oriti '14; Benedetti, Lahoche '15...]



Work within a truncation e.g. φ^6 truncation in $d = 3$:

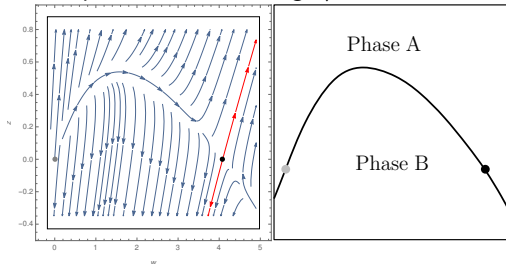
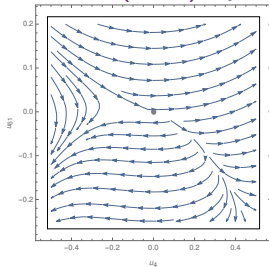
$$\begin{aligned}
 \Gamma_k = & -Z(k) \sum_{\ell=1}^3 \text{[Diagram: loop with cross]} + Z(k) k^2 u_2(k) \text{[Diagram: bubble]} + Z(k)^2 k \frac{u_4(k)}{2} \sum_{\ell=1}^3 \text{[Diagram: square]} \\
 & + Z(k)^3 \frac{u_{6,1}(k)}{3} \sum_{\ell=1}^3 \text{[Diagram: pentagon]} + Z(k)^3 u_{6,2}(k) \sum_{\ell=1}^3 \text{[Diagram: hexagon]}
 \end{aligned}$$

Wetterich (exact) equation \rightarrow truncated non-perturbative flows e.g. φ^6 :



\Rightarrow Non-perturbative UV fixed point with 1 IR-relevant direction.

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This statement is **robust under refinement** of the truncations:

$d = 3$

n	4	6	8	10	12
θ_1	2.5	2.7	2.8	2.9	3.0
θ_2	-0.37	-0.31	-0.28	-0.28	-0.31
θ_3	-	-1.7	-1.6	-1.6	-1.7
θ_4	-	-	-4.0	-4.1	-4.3
θ_5	-	-	-	-6.6	-6.7
θ_6	-	-	-	-	-9.5
η	-0.70	-0.82	-0.94	-1.1	-1.2

$d = 4$

n	4	6	8	10	12
θ_1	2.1	2.2	2.3	2.4	2.5
θ_2	-0.42	-0.33	-0.25	-0.16	-0.11
θ_3	-	-1.5	-1.3	-1.3	-1.3
θ_4	-	-	-3.4	-3.2	-3.4
θ_5	-	-	-	-5.5	-5.6
θ_6	-	-	-	-	-8.0
η	-0.64	-0.71	-0.77	-0.84	-0.95

- 1 Renormalization of spin foams: why and how?
- 2 A non-trivial but tractable GFT theory space
- 3 Perturbative and non-perturbative results
- 4 Conclusion and outlook

- GFT = QFT **completion of Spin Foam models**.
- **Colored trace invariants** \rightarrow **infinite-dimensional** but manageable GFT theory space (stability under the RG flow).
- Appreciable number of **renormalizable GFTs** known e.g. in the theory space of $SU(2)$ **gauge invariant GFTs**.
- Hints of **non-trivial fixed points** \rightarrow **phase transitions** in quantum gravity?

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Take-home message: Spin foam renormalization schemes can be defined; they rely on a delicate interplay between the **combinatorics**, the **topology** and the **geometric data** of the foams. Extra care will be required in QG models.

GFT at Loops:

Chirco, de Cesare, Finocchiaro, Freidel, Gielen, Kegeles, Kotecha, Livine, Oriti, Perez-Sanchez, Pithis, Raasakka, Rovelli, Thürigen, Wilson-Ewing, Zhang...

Not at Loops, but contributed significantly to GFT/TFT renormalization:

Baratin, Benedetti, Ben Geloun, Bonzom, Eichhorn, Gurau, Krajewski, Koslowski, Lahoche, Ousmane Samary, Rivasseau, Ryan, Smerlak, Tanasa, Toriumi, Vignes-Tourneret...

Long-term questions:

- Is there (at least) one phase in which (a sector of) continuous **general relativity** is recovered in the classical limit of some 4d Lorentzian model?
- Can these questions be addressed within a **well-controlled approximation scheme**, allowing to subsequently compute **quantum corrections** to classical general relativity?

Steps towards quantum gravity applications

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Some technical and conceptual stumbling blocks:

- **Lorentzian signature** \Rightarrow non-compact groups? [Ben Geloun, Martini, Oriti '15 '16]
- Structure of the divergences in the presence of **simplicity constraints**? [Riello '13; Bonzom, Dittrich '13; Chen '16]
- expansion around **non-trivial vacua**, and connection to **GFT condensates**? [Plenaries by Geiller and Gielen]
- **Propagators** in 4d gravity? Appropriate notion of **scale**? Role of the **Barbero-Immirzi** parameter? [Charles, Livine '15]
- **Observables** and **symmetries**? [Baratin, Girelli, Oriti '11; Baratin, Freidel, Gurau '14; Kegeles, Oriti '16; Dittrich '10s]

Strategy I: build up quantum geometries from elementary building blocks.

- **Loop quantum gravity and spin foams:**

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

→
quantization

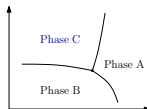


Planck-scale quantum geometry

- **GFT formulation:**

$$\sum_{\Delta} w \left(\frac{\text{Diagram}}{\Delta} \right)$$

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Effective phase diagram

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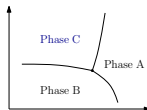


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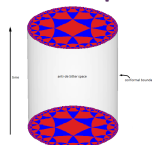


Effective phase diagram

Strategy II: Assume some weak background structure, e.g. at a boundary.

- **String theory and AdS/CFT correspondence:**

(<https://commons.wikimedia.org/>)



Bulk geometry from boundary CFT dynamics

- SYK model: N Majorana fermions ψ_i with random q -valent interactions: [Sachdev, Ye '93; Sachdev '10; Kitaev '15...]

$$H_{\text{int}} = \sum_{1 \leq i_1 < \dots < i_q \leq N} j_{i_1 \dots i_q} \psi_{i_1} \psi_{i_2} \dots \psi_{i_q}$$

$$\langle j_{i_1 \dots i_q}^2 \rangle = \frac{J^2 (q-1)!}{N^{q-1}}$$

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⇒ tractable toy-model of AdS/CFT

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[Kitaev '15; Maldacena, Stanford '16...]

- Same type of Schwinger-Dyson equations as in tensor models and GFTs ('melonic' equations) ⇒ SYK conformal regime from tensor models.

[Witten '16; Gurau '16; Klebanov, Tarnopolsky '16...]

Technical convergence of strings and loops in simplified models → unexpected dialogue opportunity!
[See talk by Verlinde, and also Pranzetti, Chirco...]

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