

Loop Equations for Gauge/Gravity Duality

JAMES SULLY

B. Czech, L. Lamprou, S.McCandlish, B. Mosk, JS arXiv: 1604.03110
B. Czech, L. Lamprou, S.McCandlish, B. Mosk, JS arXiv: 1608.06282
B. Czech, L. Lamprou, S.McCandlish, JS Ongoing

Loops 2017 – University of Warsaw July, 2017

• A well defined, UV complete, background-independent description of a gravitational theory

- A well defined, UV complete, background-independent description of a gravitational theory
 - Gravitational theory is described by an ordinary conformal field theory.

- A well defined, UV complete, background-independent description of a gravitational theory
 - Gravitational theory is described by an ordinary conformal field theory.

• Naturally (necessarily?) unified description of gravity and matter

- A well defined, UV complete, background-independent description of a gravitational theory
 - Gravitational theory is described by an ordinary conformal field theory.

- Naturally (necessarily?) unified description of gravity and matter
 - CFT spectrum contains operators that describe both matter and coupling to gravity.

- A well defined, UV complete, background-independent description of a gravitational theory
 - Gravitational theory is described by an ordinary conformal field theory.

- Naturally (necessarily?) unified description of gravity and matter
 - CFT spectrum contains operators that describe both matter and coupling to gravity.

• Most profound feature of quantum gravity is built-in: holographic

- A well defined, UV complete, background-independent description of a gravitational theory
 - Gravitational theory is described by an ordinary conformal field theory.

- Naturally (necessarily?) unified description of gravity and matter
 - CFT spectrum contains operators that describe both matter and coupling to gravity.

- Most profound feature of quantum gravity is built-in: holographic
 - The CFT lives in **one fewer dimension** than the gravity dual

• Black hole thermodynamics:

- Black hole thermodynamics:
 - Black holes are thermal states in the dual CFT and have a real temperature

- Black hole thermodynamics:
 - Black holes are thermal states in the dual CFT and have a real temperature
 - Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)

- Black hole thermodynamics:
 - Black holes are thermal states in the dual CFT and have a real temperature
 - Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
 - Black hole evaporation is a unitary process

- Black hole thermodynamics:
 - Black holes are thermal states in the dual CFT and have a real temperature
 - Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
 - Black hole evaporation is a unitary process
- Area and entropy:

• Black hole thermodynamics:

- Black holes are thermal states in the dual CFT and have a real temperature
- Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
- Black hole evaporation is a unitary process

• Area and entropy:

 More generally, certain areas in spacetime related to entanglement entropy of regions in the CFT (Hubeny-Rangamani-Ryu-Takayanagi)

• Black hole thermodynamics:

- Black holes are thermal states in the dual CFT and have a real temperature
- Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
- Black hole evaporation is a unitary process

• Area and entropy:

- More generally, certain areas in spacetime related to entanglement entropy of regions in the CFT (Hubeny-Rangamani-Ryu-Takayanagi)
- String theory... or not:

• Black hole thermodynamics:

- Black holes are thermal states in the dual CFT and have a real temperature
- Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
- Black hole evaporation is a unitary process

• Area and entropy:

- More generally, certain areas in spacetime related to entanglement entropy of regions in the CFT (Hubeny-Rangamani-Ryu-Takayanagi)
- String theory... or not:
 - String theory is a well-defined, UV complete, background-ind. theory.

• Black hole thermodynamics:

- Black holes are thermal states in the dual CFT and have a real temperature
- Black hole entropy is just the degeneracy of states at large conformal dimension in the CFT (in 2D this is just the Cardy growth of states)
- Black hole evaporation is a unitary process

• Area and entropy:

• More generally, certain **areas in spacetime** related to **entanglement entropy** of regions in the CFT (Hubeny-Rangamani-Ryu-Takayanagi)

• String theory... or not:

- **String theory** is a **well-defined**, UV complete, background-ind. theory.
- But gauge/gravity duality is more general: dual theory need not be a string theory. Expect consistent theories of gravity from CFTs that have no string theory description.

How do we describe the interior of a black hole?

How do we describe the interior of a black hole?

How does gravitational geometry *emerge* from the CFT?

How do we describe the interior of a black hole?

How does gravitational geometry *emerge* from the CFT?

... what has gauge/gravity duality taught us?

How do we describe the interior of a black hole? How does gravitational geometry *emerge* from the CFT?

... what has gauge/gravity duality taught us?

Not nearly as much as we would have hoped by now!

How do we describe the interior of a black hole? How does gravitational geometry *emerge* from the CFT?

... what has gauge/gravity duality taught us?

Not nearly as much as we would have hoped by now!

Despite being gifted with a fantastic, UV-complete theory of gravity...

How do we describe the interior of a black hole? How does gravitational geometry *emerge* from the CFT?

... what has gauge/gravity duality taught us?

Not nearly as much as we would have hoped by now!

Despite being gifted with a fantastic, UV-complete theory of gravity... ...we are still much better at understanding perturbative, background-dependent questions in low-energy effective field theory.

How do we describe the interior of a black hole?

How does gravitational geometry *emerge* from the CFT?

... what has gauge/gravity duality taught us?

Not nearly as much as we would have hoped by now!

Despite being gifted with a fantastic, UV-complete theory of gravity...

...we are still much better at understanding perturbative, background-dependent questions in low-energy effective field theory.

Problem: Less that we can't compute in the CFT, more that we don't know what to ask.

Would like to know:

Would like to know:

1. What lessons and tools can 'string theorists' borrow from other communities that have also thought hard about these problems?

Would like to know:

- 1. What lessons and tools can 'string theorists' borrow from other communities that have also thought hard about these problems?
- 2. Can experts of other approaches make progress by working within the framework of gauge/gravity duality?

Would like to know:

- 1. What lessons and tools can 'string theorists' borrow from other communities that have also thought hard about these problems?
- 2. Can experts of other approaches make progress by working within the framework of gauge/gravity duality?

This talk: Some work that leans in this direction...

AdS/CFT Duality
Conformal Field theory d-dimensions

Conformal Field theory d-dimensions **Theory of Gravity** (d+1)-dimensions

Conformal Field theory d-dimensions

Theory of Gravity (d+1)-dimensions







Conformal Field theory Theory of Gravity (d+1)-dimensions d-dimensions

Conformal Field theory d-dimensions

 \mathcal{O}_Δ

 Z_{CFT}



Conformal Field theory Theory of Gravity (d+1)-dimensions d-dimensions Z_{CFT} Z_{Grav} **Semiclassical** Limit **Classical geometry**, **Generalized Free CFT Perturbative gravity**



[cf. Heemskerk, Penedones, Polchinski, JS]

Only particular class of CFTs expected to have a 'good' gravitational dual:

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

A 'good' CFT must have:

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1
- 2. Correlators that factorize:

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1
- 2. Correlators that factorize:

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1
- 2. Correlators that factorize:

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1
- 2. Correlators that factorize:

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

$$\mathcal{O} \quad \mathcal{O}^2 \quad \mathcal{O}\partial\mathcal{O} \qquad O(c)$$

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: c >> 1
- 2. Correlators that factorize:

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

$$\mathcal{O} \quad \mathcal{O}^2 \quad \mathcal{O}\partial\mathcal{O} \qquad O(c)$$

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: $c >> 1\left(\frac{L}{l_n}\right) \gg 1$
- 2. Correlators that factorize:

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

$$\mathcal{O} \quad \mathcal{O}^2 \quad \mathcal{O}\partial\mathcal{O} \qquad O(c)$$

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: $c >> 1\left(\frac{L}{L}\right) \gg 1$
- 2. Correlators that factorize:

Perturbative effective fields in bulk

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: $c >> 1\left(\frac{L}{L}\right) \gg 1$
- 2. Correlators that factorize:

Perturbative effective fields in bulk

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A **spectrum of conformal dimensions** that is **sparse**:

Low-energy fields $O O^2 O \partial O O$



Only particular class of CFTs expected to have a 'good' gravitational dual: What are these CFTs?

- A 'good' CFT must have:
- 1. A large central charge: $c >> 1\left(\frac{L}{L}\right) \gg 1$
- 2. Correlators that factorize:

Perturbative effective fields in bulk

 $\left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle = \left\langle \mathcal{O}_1(x_1)\mathcal{O}_1(x_2)\right\rangle \left\langle \mathcal{O}_2(x_3)\mathcal{O}_2(x_4)\right\rangle + O(1/c)$

1. A spectrum of conformal dimensions that is sparse: Low-energy fields Black hole states

 $\begin{array}{c|c} & \bullet & \bullet \\ & \mathcal{O} & \mathcal{O}^2 & \mathcal{O}\partial\mathcal{O} & O(c) \end{array}$ Then this CFT is dual to a theory of gravity whose low-energy energy description is gravity plus effective field theory of field ϕ dual to \mathcal{O} .





it also describes some geometry in the gravitational theory.













Different Languages

 \sim

Different Languages

Standard object to work with in the **gauge theory** are **local operators**

Different Languages

Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.
Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



Standard object to work with **in gravity** is a **local operator** in the bulk.

Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



Standard object to work with **in gravity** is a **local operator** in the bulk.

• Gravity has **no local operators** (diff-invariant)

Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



Standard object to work with **in gravity** is a **local operator** in the bulk.

- Gravity has **no local operators** (diff-invariant)
- Quasi-local operator is complicated to specify in terms of boundary conditions on the cylinder.

Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



- Gravity has **no local operators** (diff-invariant)
- Quasi-local operator is complicated to specify in terms of boundary conditions on the cylinder.





Standard object to work with in the **gauge theory** are **local operators**

• These are UV operators that don't naturally see far into the bulk.



Standard object to work with **in gravity** is a **local operator** in the bulk.

- Gravity has **no local operators** (diff-invariant)
- Quasi-local operator is complicated to specify in terms of boundary conditions on the cylinder.



 \cdot (x)

[cf. Hijano, Kraus, Perlmutter, Snively]

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:

Gravity:

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:

Gravity:

$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle \equiv$$

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



Dynamical Parameters

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



Dynamical Parameters Conformal Kinematics: All integrals hidden here!

[cf. Hijano, Kraus, Perlmutter, Snively]

Using these non-invariant effective observables is **meant to be most straight-forward** way to do bulk calculations:



Dynamical Parameters Conformal Kinematics: All integrals hidden here!

... is being a bit more sophisticated worth the effort?

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

- 1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?
- 2. Can we find a **better connection to the natural CFT variables** that make the OPE so simple?

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

- 1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?
- 2. Can we find a **better connection to the natural CFT variables** that make the OPE so simple?

Answer:

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

- 1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?
- 2. Can we find a **better connection to the natural CFT variables** that make the OPE so simple?

Answer:

Stereoscopic Dictionary:

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

- 1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?
- 2. Can we find a **better connection to the natural CFT variables** that make the OPE so simple?

Answer:

Stereoscopic Dictionary:

Surface Operators in AdS (Non-local, diff-invariant)

... is being a bit more sophisticated worth the effort? Aim of this talk: YES.

Questions:

- 1. What is an **improved gauge/gravity dictionary** that is **gauge-invariant** on one side and **manifestly diff-invariant** on the other?
- 2. Can we find a **better connection to the natural CFT variables** that make the OPE so simple?

Answer:

Stereoscopic Dictionary:

Surface Operators in AdS (Non-local, diff-invariant)



Partial Waves of the OPE (building blocks of the OPE)

You might be familiar with a similar story:

You might be familiar with a similar story:

We typically formulate a gauge theory in terms of a **gauge potential** $A^a_\mu(X)$ and the corresponding **field strength** $F^a_{\mu\nu}$, where the equation of motion is

$$(\nabla^{\mu}F_{\mu\nu})^{a} = 0$$

You might be familiar with a similar story:

We typically formulate a gauge theory in terms of a **gauge potential** $A^a_\mu(X)$ and the corresponding **field strength** $F^a_{\mu\nu}$, where the equation of motion is

$$(\nabla^{\mu}F_{\mu\nu})^{a} = 0$$

But the inclusion of the gauge field adds redundancy to the description.

You might be familiar with a similar story:

We typically formulate a gauge theory in terms of a **gauge potential** $A^a_\mu(X)$ and the corresponding **field strength** $F^a_{\mu\nu}$, where the equation of motion is

$$(\nabla^{\mu}F_{\mu\nu})^{a} = 0$$

But the inclusion of the gauge field adds redundancy to the description. However, we do have a complete set of gauge-invariant observables

You might be familiar with a similar story:

We typically formulate a gauge theory in terms of a **gauge potential** $A^a_\mu(X)$ and the corresponding **field strength** $F^a_{\mu\nu}$, where the equation of motion is

$$(\nabla^{\mu}F_{\mu\nu})^{a} = 0$$

But the inclusion of the gauge field adds redundancy to the description. However, we do have a complete set of gauge-invariant observables

$$\Psi(C) = \operatorname{Tr}\left[\operatorname{Pexp}(\oint_{c} A_{\mu} dx^{\mu})\right] \qquad \qquad W(C) = \langle \Psi(C) \rangle$$

You might be familiar with a similar story:

We typically formulate a gauge theory in terms of a **gauge potential** $A^a_\mu(X)$ and the corresponding **field strength** $F^a_{\mu\nu}$, where the equation of motion is

$$(\nabla^{\mu}F_{\mu\nu})^{a} = 0$$

But the inclusion of the gauge field adds redundancy to the description. However, we do have a complete set of gauge-invariant observables

$$\Psi(C) = \operatorname{Tr}\left[\operatorname{Pexp}(\oint_{c} A_{\mu} dx^{\mu})\right] \qquad \qquad W(C) = \langle \Psi(C) \rangle$$

Can we express the dynamics of the theory in terms of these variables alone? (That is, can we write an equivalent EOM in the space of loops?)
Loops

[Migdal, Makeenko; Polyakov; Halpern, Makeenko;...]

















How do we insert the equation of motion into a Wilson loop?



Loop Equation: $\Box_L W = 0$ (Classically)



How do we insert the equation of motion into a Wilson loop?



Loop Equation: $\Box_L W = 0$ (Classically)

• Can also define a field space Laplacian, \Box_F , and view Wilson loop as a 'loop transform' W of the gauge field. At large N:



How do we insert the equation of motion into a Wilson loop?



Loop Equation: $\Box_L W = 0$ (Classically)

• Can also define a field space Laplacian, \Box_F , and view Wilson loop as a 'loop transform' W of the gauge field. At large N:

 $\Box_L W[A] = -W[\Box_F A] \quad \text{Intertwining Operators}$

We want to know:

We want to know:

1. What is the right **loop space** for AdS/CFT?

We want to know:

- 1. What is the right **loop space** for AdS/CFT?
- 2. What are the loop operators and their loop equations for bulk physics?

We want to know:

- 1. What is the right **loop space** for AdS/CFT?
- 2. What are the **loop operators** and their **loop equations for bulk physics**?
- 3. What do the loop equations and loop operators look like in the gauge theory?

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

• The Ryu-Takayanagi (RT/HRT) proposal connects:

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

• The Ryu-Takayanagi (RT/HRT) proposal connects:



The **entanglement entropy** of a region A on the boundary.

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

• The Ryu-Takayanagi (RT/HRT) proposal connects:



The **entanglement entropy** of a region A on the boundary.



The **area of a minimal surface** S in the gravitational dual

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

• The Ryu-Takayanagi (RT/HRT) proposal connects:



The **entanglement entropy** of a region A on the boundary.



The **area of a minimal surface** S in the gravitational dual

• Entanglement entropy of regions in the CFT gives **non-local**, **non-perturbative** probe of spacetime geometry.

We need to understand: what is the most natural loop/phase space for gravitational physics in AdS/CFT?

Thankfully an example of the type of construction we're searching for has already been found: Ryu-Takayanagi Proposal

• The Ryu-Takayanagi (RT/HRT) proposal connects:



The **entanglement entropy** of a region A on the boundary.



The **area of a minimal surface** S in the gravitational dual

• Entanglement entropy of regions in the CFT gives **non-local**, **non-perturbative** probe of spacetime geometry.

Might think of spacetime as geometrization of this entanglement data!

Let's use these non-local entanglement probes to formulate our loop space:

Let's use these non-local entanglement probes to formulate our loop space:

• One might define the relevant **loop space** to be **all possible such surfaces** (ie. all possible such boundary regions.)

Let's use these non-local entanglement probes to formulate our loop space:

• One might define the relevant **loop space** to be **all possible such surfaces** (ie. all possible such boundary regions.)



Let's use these non-local entanglement probes to formulate our loop space:

• One might define the relevant **loop space** to be **all possible such surfaces** (ie. all possible such boundary regions.)



• Enormously redundant: infinite-dimensional space

Let's use these non-local entanglement probes to formulate our loop space:

• One might define the relevant **loop space** to be **all possible such surfaces** (ie. all possible such boundary regions.)



- Enormously redundant: infinite-dimensional space
- Define our loop space to be the space of minimal surfaces with spherical boundary conditions (in any frame).
 - We will give it a new name: Kinematic Space.

What does kinematic space look like?

What does kinematic space look like?

• Consider an ordered pair of timelike separated points on the boundary:

What does kinematic space look like?

• Consider an ordered pair of timelike separated points on the boundary:



Boundary

What does kinematic space look like?

• Consider an ordered pair of timelike separated points on the boundary:



Boundary

What does kinematic space look like?

• Consider an ordered pair of timelike separated points on the boundary:



What does kinematic space look like?

• Consider an ordered pair of timelike separated points on the boundary:


Can we assign a metric to kinematic space?

Can we assign a metric to kinematic space?

- For the AdS vacuum, the metric on ${\cal K}$ is uniquely fixed by isometries of the geometry:

$$ds^2 = \frac{I_{\mu\nu}(l^{\mu})}{l^2} \left(d\chi^{\mu} d\chi^{\nu} - dl^{\mu} dl^{\nu} \right) \qquad \qquad \mathbf{\chi}: \mathbf{Center}$$

l: Separation

Can we assign a metric to kinematic space?

- For the AdS vacuum, the metric on ${\cal K}$ is uniquely fixed by isometries of the geometry:

$$ds^{2} = \frac{I_{\mu\nu}(l^{\mu})}{l^{2}} \left(d\chi^{\mu} d\chi^{\nu} - dl^{\mu} dl^{\nu} \right)$$

Signature: (d,d)

 χ : Center

l: Separation

Can we assign a metric to kinematic space?

- For the AdS vacuum, the metric on ${\cal K}$ is uniquely fixed by isometries of the geometry:

$$ds^{2} = \frac{I_{\mu\nu}(l^{\mu})}{l^{2}} \left(d\chi^{\mu} d\chi^{\nu} - dl^{\mu} dl^{\nu} \right)$$

χ: Center

l: Separation

Signature: (d,d)

$$ds^{2} = \frac{1}{2} \left[\frac{dq^{2} - dl^{2}}{l^{2}} + \frac{d\bar{q}^{2} - d\bar{l}^{2}}{\bar{l}^{2}} \right] = \frac{1}{2} \left[ds_{z}^{2} + ds_{\bar{z}}^{2} \right]$$

Can we assign a metric to kinematic space?

- For the AdS vacuum, the metric on ${\cal K}$ is uniquely fixed by isometries of the geometry:



Can we assign a metric to kinematic space?

- For the AdS vacuum, the metric on ${\cal K}$ is uniquely fixed by isometries of the geometry:



We have a 'loop space', but still need the right 'Wilson Loops'

We have a 'loop space', but still need the right 'Wilson Loops'

• We can think of their areas as integrating the unit operator over the minimal surface:

$$A = \int d^n x \sqrt{h}(1)$$

We have a 'loop space', but still need the right 'Wilson Loops'

• We can think of their areas as integrating the unit operator over the minimal surface:

$$A = \int d^n x \sqrt{h}(1)$$



• A natural generalization then is:

$$\tilde{\phi} = \int d^n x \sqrt{h}(\phi)$$



We have a 'loop space', but still need the right 'Wilson Loops'

• We can think of their areas as integrating the unit operator over the minimal surface:

$$A = \int d^n x \sqrt{h}(1)$$



• A natural generalization then is:

$$\tilde{\phi} = \int d^n x \sqrt{h}(\phi)$$



• This bulk surface/geodesic operator is a **non-local** and **diff-invariant** bulk probe.

We have a 'loop space', but still need the right 'Wilson Loops'

• We can think of their areas as integrating the unit operator over the minimal surface:

$$A = \int d^n x \sqrt{h}(1)$$



• A natural generalization then is:

$$\tilde{\phi} = \int d^n x \sqrt{h}(\phi)$$
Radon Transform: $R[\phi] \coloneqq \widetilde{\phi}$

• This bulk surface/geodesic operator is a **non-local** and **diff-invariant** bulk probe.



















Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

$$\left(\Box_{AdS} - m^2\right)\phi(x) = 0$$

Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

$$\left(\Box_{AdS} - m^2\right)\phi(x) = 0 \quad \longleftrightarrow \quad \Box_{KS}Rf = -R\Box_{AdS}f$$

Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

$$\left(\Box_{AdS} - m^2\right)\phi(x) = 0 \quad \longleftrightarrow \quad \Box_{KS}Rf = -R\Box_{AdS}f$$
$$\left(\Box_{KS} + m^2)\tilde{\phi}(\gamma) = 0\right)$$

Intertwinement allows us to **rewrite the dynamics of the gravitational** theory in terms of **dynamics on kinematic space**:

$$\left(\Box_{AdS} - m^2 \right) \phi(x) = 0 \quad \longleftrightarrow \quad \Box_{KS} Rf = -R \Box_{AdS} f$$

$$\left(\Box_{KS} + m^2) \tilde{\phi}(\gamma) = 0 \right) \quad \mathsf{Loop EOM for free scalar}$$

$$= \mathsf{Kinematic Free scalar EOM}$$

We haven't succeeded unless these appear naturally in the dual gauge theory:

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

 $\mathcal{O}_i(x)$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

 $\mathcal{O}_{i}(x)$ $\mathcal{O}_{i}(0)$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:


We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

$$\mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(0) \\ * \\ \mathcal{O}_{i}(0) \\ * \\ \mathcal{O}_{i}(0) \\ \mathcal{O}_{i}(x) \mathcal{O}_{i}(0) = \sum_{k} C_{iik} |x|^{\Delta_{k} - 2\Delta_{i}} (1 + b_{1} x^{\mu} \partial_{\mu} + b_{2} x^{\mu} x^{\nu} \partial_{\mu} \partial_{\nu} + \dots) \mathcal{O}_{k}(0)$$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

$$\mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(0) \\ \mathbb{X} \\ \mathcal{O}_{i}(0) \\ \mathbb{X} \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(0) = \sum_{k} C_{iik} |x|^{\Delta_{k} - 2\Delta_{i}} (1 + b_{1} x^{\mu} \partial_{\mu} + b_{2} x^{\mu} x^{\nu} \partial_{\mu} \partial_{\nu} + \dots) \\ \mathcal{O}_{k}(0) \\ \mathbb{C}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal$$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:

$$\mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(0) \\ \mathbb{D}_{i}(0) \\ \mathbb{D}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(x) \\ \mathcal{O}_{i}(0) = \sum_{k} C_{iik} |x|^{\Delta_{k} - 2\Delta_{i}} (1 + b_{1} x^{\mu} \partial_{\mu} + b_{2} x^{\mu} x^{\nu} \partial_{\mu} \partial_{\nu} + \dots) \\ \mathcal{O}_{k}(0) \\ \mathbb{D}_{i}(x) \\ \mathbb$$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:



• Let us introduce a more compact notation for this expansion

$$\mathcal{O}_{i}(x) \mathcal{O}_{i}(y) = |x - y|^{-2\Delta_{i}} \sum_{k} C_{iik} \mathcal{B}_{k}(x, y)$$

We haven't succeeded unless these appear naturally in the dual gauge theory:

• Consider a quasi-primary operator $\mathcal{O}_i(x)$ of dimensions Δ_i . We can expand the product of two such operators using a local basis of operators:



• Let us introduce a more compact notation for this expansion

$$\mathcal{O}_{i}(x) \mathcal{O}_{i}(y) = |x - y|^{-2\Delta_{i}} \sum_{k} C_{iik} \mathcal{B}_{k}(x, y)$$

• We will call $\mathcal{B}_k(x,y)$ the 'OPE Block'

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

Consider a scalar block $(\Delta_k, l = 0)$.

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

Consider a scalar block $(\Delta_k, l = 0)$. What is its equation of motion?

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

Consider a scalar block $(\Delta_k, l = 0)$.

What is its equation of motion?

• Eigen-operator of the conformal Casimir: $[L^2, \mathcal{B}_k(x, y)] = C_{\mathcal{O}_k} \mathcal{B}_k(x, y)$

 $C_{\mathcal{O}_k} = -\Delta \left(\Delta - d \right)$

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

 $C_{\mathcal{O}_k} = -\Delta \left(\Delta - d \right)$

Consider a scalar block $(\Delta_k, l = 0)$.

What is its equation of motion?

- Eigen-operator of the conformal Casimir: $[L^2, \mathcal{B}_k(x, y)] = C_{\mathcal{O}_k} \mathcal{B}_k(x, y)$
- We represent this as

$$\mathcal{L}^2_{(B)} = 2\Box_{\mathrm{KS}}$$

• The OPE block carries coordinates of two points (x, y), so we might naturally identify it with a **field living in kinematic space**.

Consider a scalar block $(\Delta_k, l = 0)$.

What is its equation of motion?

- Eigen-operator of the conformal Casimir: $[L^2, \mathcal{B}_k(x, y)] = C_{\mathcal{O}_k} \mathcal{B}_k(x, y)$
- We represent this as

$$\mathcal{L}^2_{(B)} = 2\Box_{\mathrm{KS}}$$

$$C_{\mathcal{O}_k} = -\Delta \left(\Delta - d\right)$$

$$\begin{bmatrix} \Box_{\rm KS} + m_{\Delta_k}^2 \end{bmatrix} \mathcal{B}_k(x, y) = 0$$
$$m_{\Delta_k}^2 = -C_{\Delta_k}$$

	Gravity	CFT
EOM		
BCs		

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs		

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs	$\lim_{x \to 0} \tilde{\phi}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$	$\lim_{x \to 0} \mathcal{B}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs	$\lim_{x \to 0} \tilde{\phi}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$	$\lim_{x \to 0} \mathcal{B}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$
Constraint		

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs	$\lim_{x \to 0} \tilde{\phi}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$	$\lim_{x \to 0} \mathcal{B}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$
Constraint	'John's Equations'	

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs	$\lim_{x \to 0} \tilde{\phi}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$	$\lim_{x \to 0} \mathcal{B}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$
Constraint	'John's Equations'	'Spin'

It is the same wave equation obeyed by a bulk geodesic operator dual to the operator \mathcal{O}_k :

	Gravity	CFT
EOM	$\left[\Box_{KS} + m_{\Delta_k}^2\right]\tilde{\phi_k}(\gamma) = 0$	$\left[\Box_{\mathrm{KS}} + m_{\Delta_k}^2\right] \mathcal{B}_k\left(x, y\right) = 0$
BCs	$\lim_{x \to 0} \tilde{\phi}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$	$\lim_{x \to 0} \mathcal{B}_k(x,0) = x^{\Delta_k} \mathcal{O}_k(0)$
Constraint	'John's Equations'	'Spin'

The Kinematic Dictionary:

$$\mathcal{B}_k\left(x,y
ight)=\, ilde{\phi}ig(\gammaig)$$

[cf. de Boer, Haehl, Heller, Myers; da Cunha, Guica]

Can we do carry out a similar procedure for the metric itself?

Can we do carry out a similar procedure for the metric itself?

• We have the linearized EOM: $\delta R_{\mu\nu}(x) = \delta T_{\mu\nu}(x) - \frac{1}{d-1} \delta T g_{\mu\nu} - d\delta g_{\mu\nu}$

Can we do carry out a similar procedure for the metric itself?

- We have the linearized EOM: $\delta R_{\mu\nu}(x) = \delta T_{\mu\nu}(x) \frac{1}{d-1} \delta T g_{\mu\nu} d\delta g_{\mu\nu}$
- Define Radon Transforms: $R_{\parallel}[s_{\mu\nu}] = \int_{B} \sqrt{h} h^{ab} \delta g_{ab}$ $R_{\perp}[s_{\mu\nu}] = \int_{B} \sqrt{h} (g^{ab} - h^{ab}) \delta g_{ab}$

and prove an analogous intertwining relation: $R_{||} [2\delta R_{\mu\nu}] = \Box_K R_{||} [\delta g_{\mu\nu}]$

Can we do carry out a similar procedure for the metric itself?

- We have the linearized EOM: $\delta R_{\mu\nu}(x) = \delta T_{\mu\nu}(x) \frac{1}{d-1} \delta T g_{\mu\nu} d\delta g_{\mu\nu}$
- Define Radon Transforms: $R_{\parallel}[s_{\mu\nu}] = \int_{B} \sqrt{h} h^{ab} \delta g_{ab}$ $R_{\perp}[s_{\mu\nu}] = \int_{B} \sqrt{h} (g^{ab} - h^{ab}) \delta g_{ab}$

and prove an analogous intertwining relation: $R_{||} [2\delta]$

 $R_{||} \left[2\delta R_{\mu\nu} \right] = \Box_K R_{||} \left[\delta g_{\mu\nu} \right]$

• We find the **equivalent Kinematic EOM**:

Can we do carry out a similar procedure for the metric itself?

- We have the linearized EOM: $\delta R_{\mu\nu}(x) = \delta T_{\mu\nu}(x) \frac{1}{d-1} \delta T g_{\mu\nu} d\delta g_{\mu\nu}$
- Define Radon Transforms: $R_{\parallel}[s_{\mu\nu}] = \int_{B} \sqrt{h} h^{ab} \delta g_{ab}$ $R_{\perp}[s_{\mu\nu}] = \int_{B} \sqrt{h} (g^{ab} - h^{ab}) \delta g_{ab}$

and prove an analogous **intertwining relation**:

 $R_{||} \left[2\delta R_{\mu\nu} \right] = \Box_K R_{||} \left[\delta g_{\mu\nu} \right]$

• We find the equivalent Kinematic EOM:

$$\left(\Box_{KS} + 2d\right) R_{\parallel} \left[\delta g\right] = -2R_{\perp} \left[\delta T\right]$$

Cf. [de Boer, Heller, Myers, Neiman] [Nozaki, Numasawa, Prudenziati, Takayanagi], [Bhattacharya, Takayanagi]

Can we do carry out a similar procedure for the metric itself?

- We have the linearized EOM: $\delta R_{\mu\nu}(x) = \delta T_{\mu\nu}(x) \frac{1}{d-1} \delta T g_{\mu\nu} d\delta g_{\mu\nu}$
- Define Radon Transforms: $R_{\parallel}[s_{\mu\nu}] = \int_{B} \sqrt{h} h^{ab} \delta g_{ab}$ $R_{\perp}[s_{\mu\nu}] = \int_{B} \sqrt{h} (g^{ab} - h^{ab}) \delta g_{ab}$

and prove an analogous intertwining relation:

 $R_{||} \left[2\delta R_{\mu\nu} \right] = \Box_K R_{||} \left[\delta g_{\mu\nu} \right]$

• We find the equivalent Kinematic EOM:

$$\left(\Box_{KS} + 2d\right) R_{\parallel} \left[\delta g\right] = -2R_{\perp} \left[\delta T\right]$$

Cf. [de Boer, Heller, Myers, Neiman] [Nozaki, Numasawa, Prudenziati, Takayanagi], [Bhattacharya, Takayanagi]

• CFT derivation: follows naturally from the Casimir equation for the stress tensor.

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

 Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

- Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.
- Bulk locality can be understood as a CFT concept:

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

- Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.
- Bulk locality can be understood as a CFT concept:



[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

 Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.

Bulk locality can be understood as a CFT concept:



[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

 Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.

Bulk locality can be understood as a CFT concept:



[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

 Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.

Bulk locality can be understood as a CFT concept:



Find a tower of corrections that restore locality and determine bulk EOM:
Local Bulk from the CFT

[Czech, Lamprou, McCandlish, JS: Ongoing; cf. Kabat et al.]

Much of what has been discussed so far is fixed by kinematics. Does this extend to dynamics in the bulk?

 Would like to do extend using the CFT alone—to see how quasi-local gravitational physics *emerges*.

Bulk locality can be understood as a CFT concept:



Find a tower of corrections that restore locality and determine bulk EOM:

$$\tilde{\phi}_{\Delta}(\gamma) = \mathcal{B}_{\Delta}(\gamma) + \frac{1}{c} \sum_{n} a_n^{CFT} \mathcal{B}_n(\gamma) + \dots$$

It's hard to specify quasi-local operators in a background-independent manner.

It's hard to specify quasi-local operators in a background-independent manner.

It's hard to specify quasi-local operators in a background-independent manner.

$$\tilde{\phi}(\gamma) = \int_{C_{\gamma}} dx \ \phi(x)$$

It's hard to specify quasi-local operators in a background-independent manner.

$$\tilde{\phi}(\gamma) = \int_{C_{\gamma}} dx \ \phi(x) \qquad \phi(x) \sim \int_{\Gamma_x} d\gamma \ \tilde{\phi}(\gamma)$$

It's hard to specify quasi-local operators in a background-independent manner.



It's hard to specify quasi-local operators in a background-independent manner.

In a fixed background: $\tilde{\phi}(\gamma) = \int_{C_{\gamma}} dx \ \phi(x) \qquad \phi(x) \sim \int_{\Gamma_{x}} d\gamma \ \tilde{\phi}(\gamma)$ In some new state:



It's hard to specify quasi-local operators in a background-independent manner.



In some new state:



Geodesic operator simply defined in terms of boundary conditions in any state—background independent

Can the same operator describe geodesics in different states?

Can the same operator describe geodesics in different states?

Can the same operator describe geodesics in different states?

Can the same operator describe geodesics in different states?

$$\langle \tilde{\phi} \dots \rangle \sim (1 + \langle T \rangle + \dots) \langle \mathcal{O} \dots \rangle$$

Can the same operator describe geodesics in different states?



Can the same operator describe geodesics in different states?



Can the same operator describe $ge \rightarrow resummation of perturbative expansion.$

The interacting block classicalizes and becomes a simple functional of the expectation values at leading order:



Resummation of operators in OPE \rightarrow resummation of perturbative expansion.

Gauge/gravity duality is a **powerful**, **UV complete model** for testing ideas about quantum gravity.

Gauge/gravity duality is a **powerful**, **UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

Gauge/gravity duality is a **powerful, UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

We found a set of **non-local**, **diff/gauge-invariant building blocks** both in the bulk and on the boundary to build a **'better' holographic dictionary**.

Gauge/gravity duality is a **powerful, UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

We found a set of **non-local**, **diff/gauge-invariant building blocks** both in the bulk and on the boundary to build a **'better' holographic dictionary**.

• On the boundary: **OPE blocks**

Gauge/gravity duality is a **powerful, UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

We found a set of **non-local**, **diff/gauge-invariant building blocks** both in the bulk and on the boundary to build a **'better' holographic dictionary**.

- On the boundary: **OPE blocks**
- In the bulk: Geodesic/Surface Operators

Gauge/gravity duality is a **powerful, UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

We found a set of **non-local**, **diff/gauge-invariant building blocks** both in the bulk and on the boundary to build a **'better' holographic dictionary**.

- On the boundary: **OPE blocks**
- In the bulk: Geodesic/Surface Operators

The simple kinematic correspondence can be extended to describe a **fullyinteracting quasi-local theory** and give **background-independent operators**.

Gauge/gravity duality is a **powerful, UV complete model** for testing ideas about quantum gravity.

But the difficult (and most interesting!) problems in QG still **require new tools to make progress**.

We found a set of **non-local**, **diff/gauge-invariant building blocks** both in the bulk and on the boundary to build a **'better' holographic dictionary**.

- On the boundary: **OPE blocks**
- In the bulk: Geodesic/Surface Operators

The simple kinematic correspondence can be extended to describe a **fullyinteracting quasi-local theory** and give **background-independent operators**.

Can we do even better by borrowing tools/ideas/people from LQG?